

Perspective of extended Higgs sectors in Super-Symmetric beyond Standard Model scenarios part-II

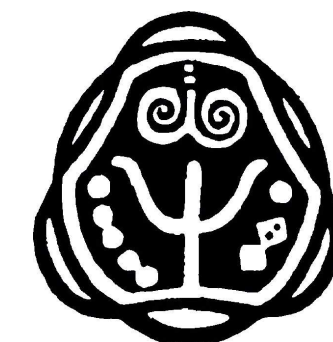
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Plan

- TMSSM: Triplino and singlino
- TNMSSM: Possible displaced vertex
- MSSM: Higgs sector CP-violation
- NMSSM: Spontaneous CP-violation

Neutralinos and charginos

- The extension of Higgs sector increases both neutralino and chargino sectors

- Neutralinos: $\tilde{B}, \tilde{W}_3, \tilde{H}_u, \tilde{H}_d, \tilde{T}_0, \tilde{S}$.

Triplino
& singlino

$$M_{\tilde{\chi}^0} = \begin{pmatrix} M_1 & 0 & -\frac{1}{2}g_Y v_d & \frac{1}{2}g_Y v_u & 0 & 0 \\ 0 & M_2 & \frac{1}{2}g_L v_d & -\frac{1}{2}g_L v_u & 0 & 0 \\ -\frac{1}{2}g_Y v_d & \frac{1}{2}g_L v_d & 0 & \frac{1}{2}v_T \lambda_T - \frac{1}{\sqrt{2}}v_S \lambda_S & \frac{1}{2}v_u \lambda_T & -\frac{1}{\sqrt{2}}v_u \lambda_S \\ \frac{1}{2}g_Y v_u & -\frac{1}{2}g_L v_u & \frac{1}{2}v_T \lambda_T - \frac{1}{\sqrt{2}}v_S \lambda_S & 0 & \frac{1}{2}v_d \lambda_T & -\frac{1}{\sqrt{2}}v_d \lambda_S \\ 0 & 0 & \frac{1}{2}v_u \lambda_T & \frac{1}{2}v_d \lambda_T & \sqrt{2}v_S \lambda_{TS} & \sqrt{2}v_T \lambda_{TS} \\ 0 & 0 & -\frac{1}{\sqrt{2}}v_u \lambda_S & -\frac{1}{\sqrt{2}}v_d \lambda_S & \sqrt{2}v_T \lambda_{TS} & \sqrt{2}\kappa v_S. \end{pmatrix}$$

- Charginos: $\tilde{W}^+, \tilde{H}_u^+, \tilde{T}_2^+ (\tilde{W}^-, \tilde{H}_d^-, \tilde{T}_1^-)$

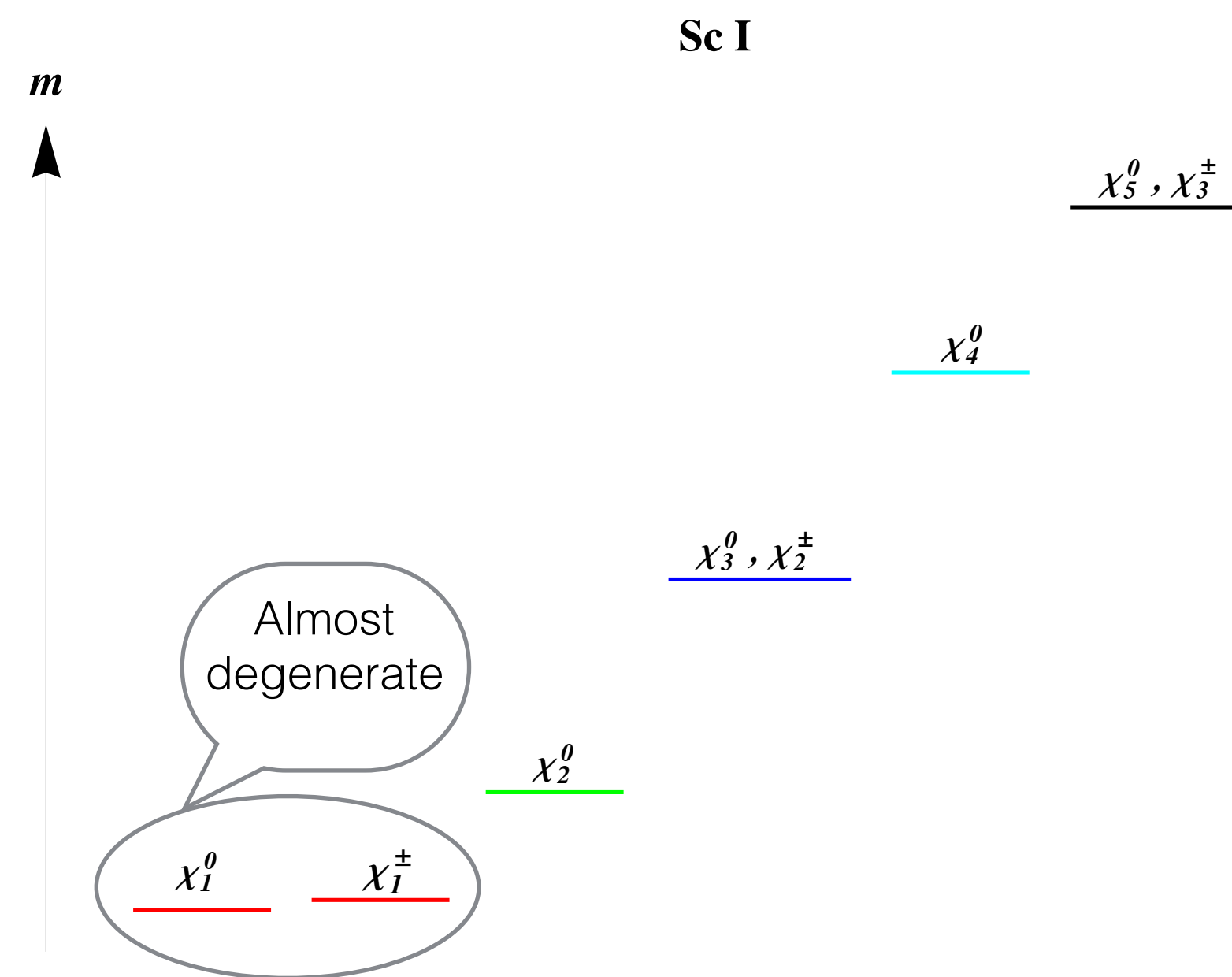
Triplino

$$M_{\tilde{\chi}^\pm} = \begin{pmatrix} M_2 & \frac{1}{\sqrt{2}}g_L v_u & -g_L v_T \\ \frac{1}{\sqrt{2}}g_L v_d & \frac{1}{\sqrt{2}}v_S \lambda_S + \frac{1}{2}v_T \lambda_T & \frac{1}{\sqrt{2}}v_u \lambda_T \\ g_L v_T & -\frac{1}{\sqrt{2}}v_d \lambda_T & \sqrt{2}v_S \lambda_{TS} \end{pmatrix}$$

Weakino spectrum

- Generic spectrum is very compressed

Benchmark Points	LSP mass (GeV)	NLSP mass (GeV)	NNLSP mass (GeV)	τ_{NLSP} (ns)	$c\tau_{NLSP}$ (cm)
BP1	542.30	542.50	864.70	0.79	23.61
BP2	561.12	561.54	651.82	0.022	0.646
BP3	530.75	530.94	771.40	1.21	36.15
BP4	498.38	498.53	722.40	3.27	97.76



- For a triplino-like LSP the NLSP is triplino-like chargino
- Very soft displaced charged leptons are produced

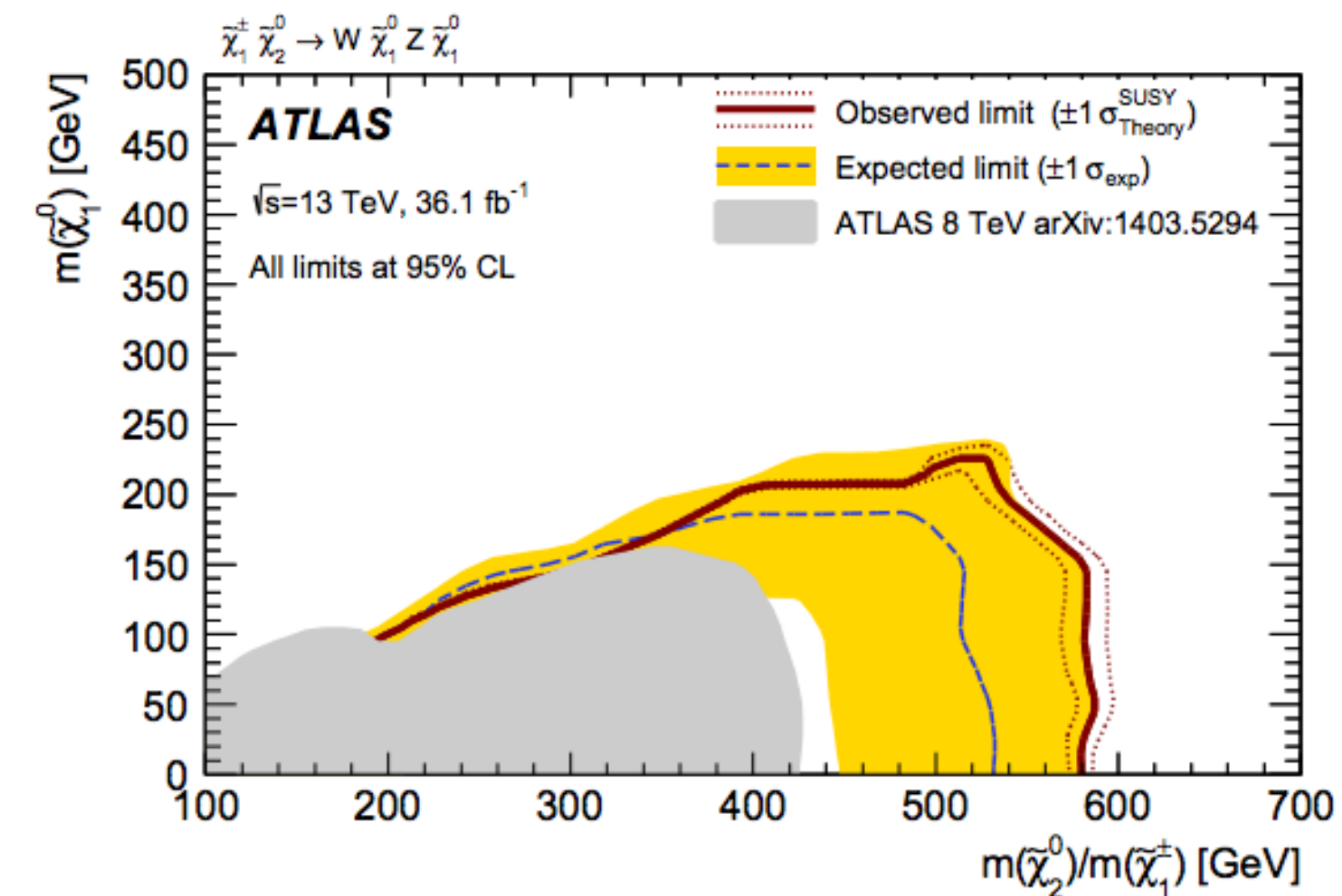
LHC Bounds from 3l , 2l signatures

- ATLAS puts a bound at 36.1 1/fb of $E_{cm}=13$ TeV masses of NLPS < 580 GeV are ruled out for massless LSP

arXiv:1803.02762 [hep-ex].

- The trigger used for these searches are with $p_T^\ell \geq 20 - 25$ GeV
- CMS also puts similar bounds

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(d) 2ℓ +jets and 3ℓ channels:
 $\tilde{\chi}_1^\pm \tilde{\chi}_2^0$ production with decays via W/Z

Triplinos with soft final states

- Triplino-type LSP decays to rather soft final states
- Often can be missed as they fall below the trigger cuts
- When the mass difference between chargino NLSP and LSP is $O(150)$ MeV,

$$\tilde{\chi}_1^\pm \rightarrow \pi^\pm \tilde{\chi}_0 \text{ mode is open and it becomes dominant}$$

- Having very low momentum NLSP can just leaves a disappearing track in the detector
- The recoil by a hard ISR-jet or the boost of the final state particle can give some visibility to the final states.

Events with ISR jet

- Hard initial state radiation (ISR) can give the hard momentum recoil to the final state
- Without ISR jet as on shell particles are back to back and degenerate with LSP $p_T' \simeq 0$
- The jets are leptons are also soft

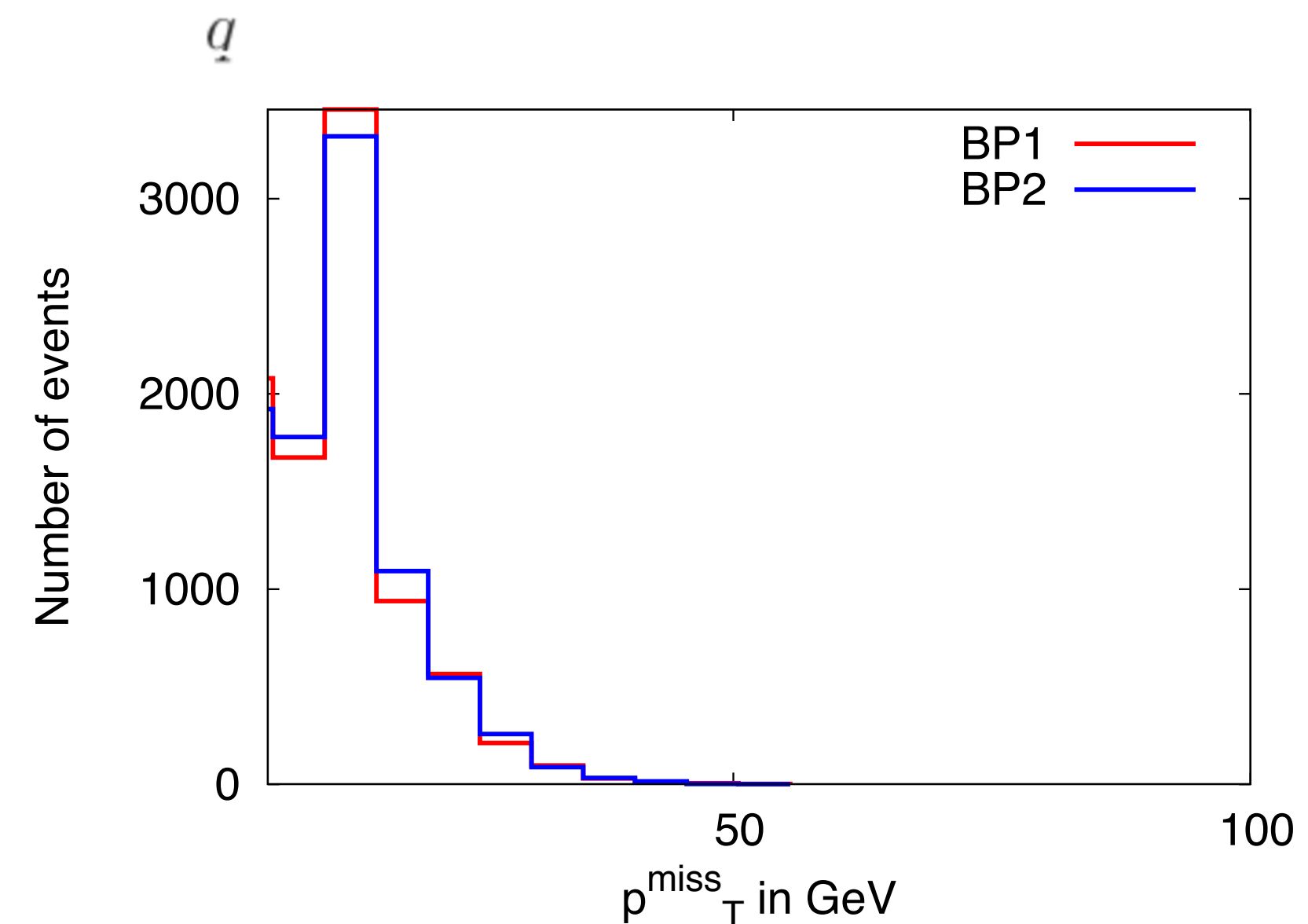
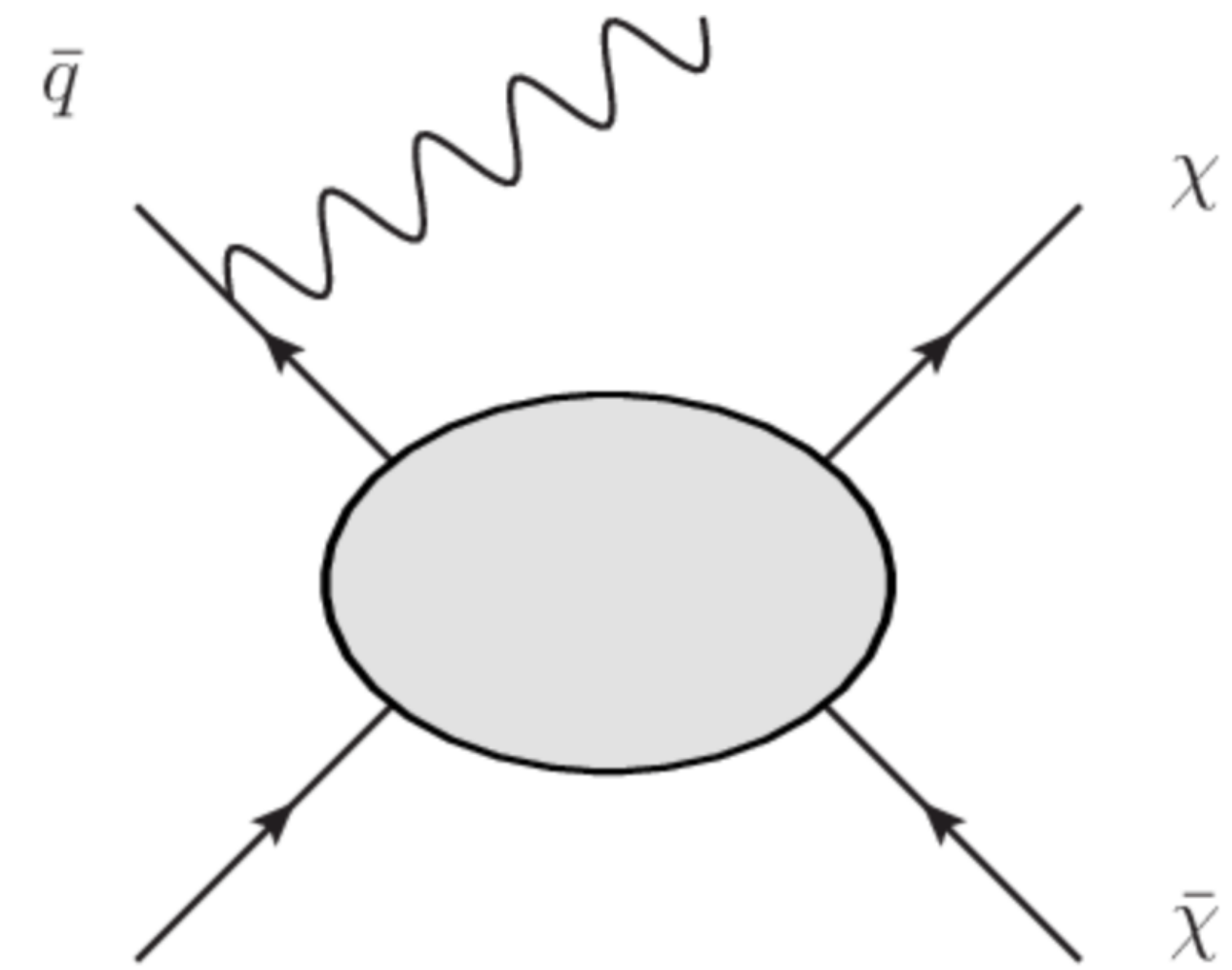
$$p_T^{\text{jets}}, p_T^{\ell} \simeq 0$$

- But with a hard ISR jet,

$$\sum p_T^i = 0$$

$$p_T^{\text{ISR}} + \underbrace{p_T^{\text{jets}} + p_T^{\ell}} + p_T' = 0$$

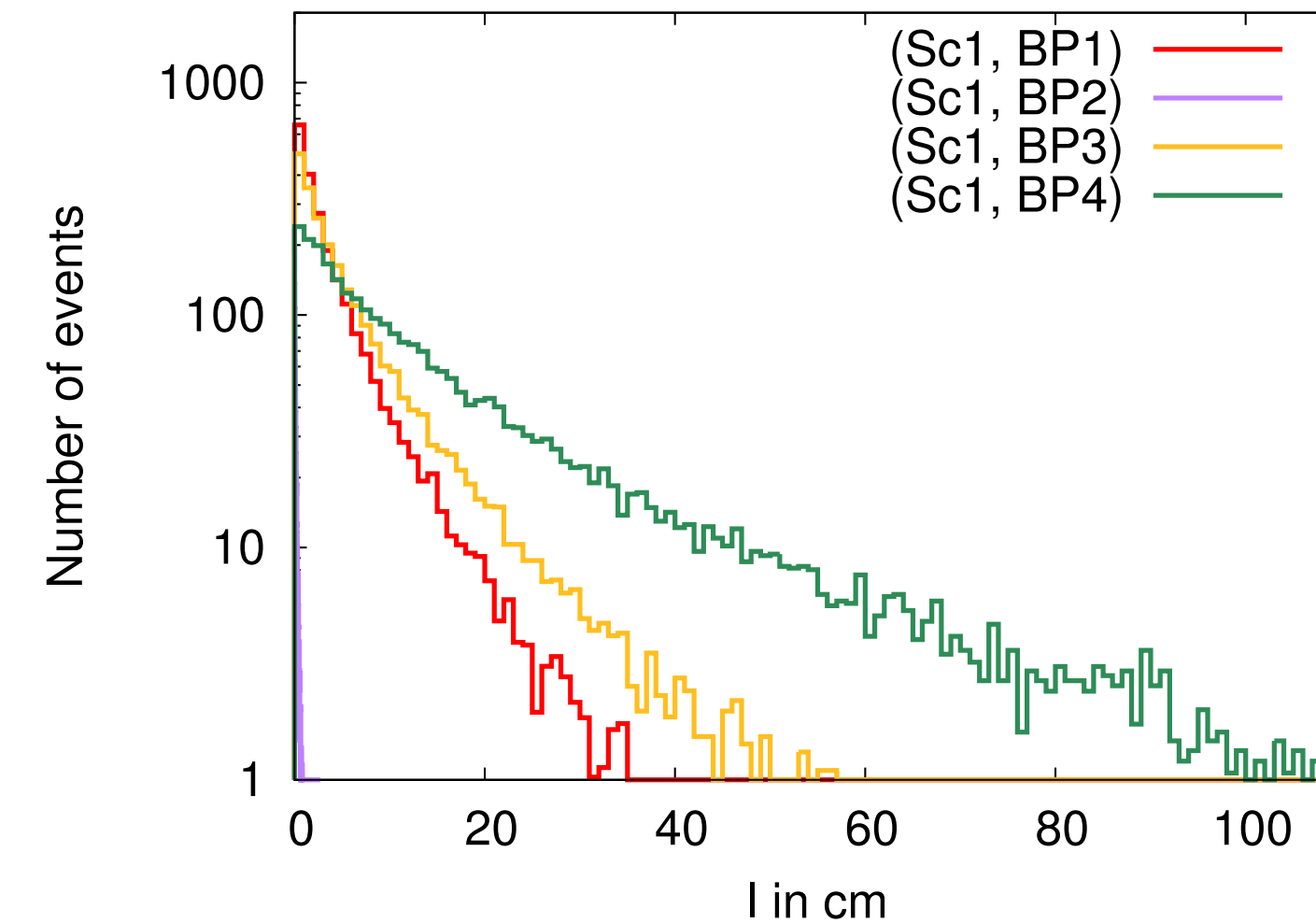
can be hard



Displaced Signature

- The displaced decay can occur from 0.1 mm to 10 m
- Displaced charged tracks and decay signatures are Standard Model background free
- $\tilde{\chi}_1^\pm \rightarrow \ell^\pm \nu \tilde{\chi}_1^0$ can give rise to mono- or di-displaced charged leptonic signatures
- $\tilde{\chi}_1^\pm \rightarrow qq' \tilde{\chi}_1^0$ gives rise to displaced jets
- For neutralino NLSP we can get di-displaced tri- and four-leptonic signatures via

$$\tilde{\chi}_2^0 \rightarrow \ell^+ \ell^- \tilde{\chi}_1^0$$



Prospect at the LHC

- Scenario 1: Triplino-type chargino NLSP- triplino LSP
- Scenario 2: Triplino-type neutralino NLSP - Bino LSP
- Scenario 3: Triplino-type chargino NLSP -Higgsino LSP
- Depending on scenarios and benchmark points these can be probed at the LHC with very early data of 30-1000 $1/\text{fb}$ with $E_{\text{cm}}=14 \text{ TeV}$
- Soft lepton and jets can be produced from long lived particle
- Soft physics is important

CP-Violating Higgs sectors

CP-Violating MSSM

- CP-violation in nature is required to explain Baryogenesis
- δ_{CKM} is the only source of CP-violation in Standard Model
- In Minimal Supersymmetric extension of Standard Model(MSSM) is a good candidate as it can have many CP-phases through various soft terms and supersymmetric μ term.
- In the model considered here, Higgs sector CP-violation enters via loop effects

Possibilities of CP-violation in MSSM

The superpotential for the MSSM is

$$W_{\text{MSSM}} = \bar{u}_i y_{\mathbf{u}i} Q_i H_u - \bar{d}_i y_{\mathbf{d}i} Q_i H_d - \bar{e}_i y_{\mathbf{e}i} L_i H_d + \mu H_u H_d .$$

- Scalar potential

$$V = (|\mu|^2 + m_{H_u}^2)|H_u^0|^2 + (|\mu|^2 + m_{H_d}^2)|H_d^0|^2 - (b H_u^0 H_d^0 + \text{c.c.}) \\ + \frac{1}{8}(g^2 + g'^2)(|H_u^0|^2 - |H_d^0|^2)^2.$$

- Soft supersymmetry-breaking terms can introduce a large number of CP-violating phases not found in the Standard

- Minimization condition: $\langle H_u \rangle = v_u, \langle v_d \rangle = v_d e^{i\theta_1}, \quad b = |b| e^{i\theta_2}$

$$\text{Im}\{b H_u H_d\} = 0 \implies \theta_1 + \theta_2 = 0$$

- No Higgs sector CP-violation at the tree-level
- In non-SUSY 2HDM we can have spontaneous CP-violation

Sources of CP violating phases in MSSM

- In the MSSM, CP-violating phases appear in the μ term of the superpotential,

$$W \supset \mu \hat{H}_u \cdot \hat{H}_d$$

- and in the soft-SUSY breaking terms as follows:

$$\begin{aligned} -\mathcal{L}_{\text{soft}} \supset & \\ & \frac{1}{2} (M_3 \tilde{g}\tilde{g} + M_2 \tilde{W}\tilde{W} + M_1 \tilde{B}\tilde{B} + \text{h.c.}) \\ & + \tilde{Q}^\dagger \mathbf{M}_{\tilde{Q}}^2 \tilde{Q} + \tilde{L}^\dagger \mathbf{M}_{\tilde{L}}^2 \tilde{L} + \tilde{u}_R^* \mathbf{M}_{\tilde{u}}^2 \tilde{u}_R + \tilde{d}_R^* \mathbf{M}_{\tilde{d}}^2 \tilde{d}_R + \tilde{e}_R^* \mathbf{M}_{\tilde{e}}^2 \tilde{e}_R \\ & - m_1^2 H_d^* H_d - m_2^2 H_u^* H_u - (m_{12}^2 H_u H_d + \text{h.c.}) \\ & + (\tilde{u}_R^* \mathbf{A}_u \tilde{Q} H_u - \tilde{d}_R^* \mathbf{A}_d \tilde{Q} H_d - \tilde{e}_R^* \mathbf{A}_e \tilde{L} H_d + \text{h.c.}) \end{aligned}$$

CP violating phases in MSSM

- Not all are independent.
- Physical observables depend on the two combinations:

$$\text{Arg}(M_i \mu (m_{12}^2)^*), \quad \text{Arg}(A_f \mu (m_{12}^2)^*),$$

with $i = 1 - 3$ and $f = e, \mu, \tau; u, c, t, d, s, b$.

- Most relevant CP phases pertinent to the Higgs sector:

$$\Phi_i \equiv \text{Arg}(M_i); \quad \Phi_{A_{f_3}} \equiv \text{Arg}(A_{f_3}),$$

with $f_3 = \tau, t, b$.

CP violating phases in MSSM

- CP violation in the Higgs potential of the MSSM leads to mixing terms between the CP-even and CP-odd Higgs fields.
[Pilaftsis, etal; 88,98](#)
- In the weak basis (G^0, a, ϕ_1, ϕ_2) , the neutral Higgs-boson mass matrix \mathcal{M}_0^2 may be cast into the form

$$\mathcal{M}_0^2 = \begin{pmatrix} \widehat{\mathcal{M}}_P^2 & \mathcal{M}_{PS}^2 \\ \mathcal{M}_{SP}^2 & \mathcal{M}_S^2 \end{pmatrix}$$

where,

$$\widehat{\mathcal{M}}_P^2 \Rightarrow \begin{pmatrix} G^0 \\ a \end{pmatrix} \leftrightarrow \begin{pmatrix} G^0 \\ a \end{pmatrix} \quad \mathcal{M}_S^2 \Rightarrow \begin{pmatrix} \phi_1 \\ \phi_2 \end{pmatrix} \leftrightarrow \begin{pmatrix} \phi_1 \\ \phi_2 \end{pmatrix}$$

$$\mathcal{M}_{PS}^2 = (\mathcal{M}_{SP}^2)^T \Rightarrow \begin{pmatrix} G^0 \\ a \end{pmatrix} \leftrightarrow \begin{pmatrix} \phi_1 \\ \phi_2 \end{pmatrix}$$

CP violating phases in MSSM

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$$\mathcal{M}_0^2 = \begin{pmatrix} \widehat{\mathcal{M}}_P^2 & \mathcal{M}_{PS}^2 \\ \mathcal{M}_{SP}^2 & \mathcal{M}_S^2 \end{pmatrix}$$

$$\Rightarrow M_{h_1} \leq M_{h_2} \leq M_{h_3} .$$

where these are not CP eigenstate.

- The mixing become significant when $(\mu A_t / M_{SUSY}^2)$ is large. and the lightest mass eigenstate becomes very light and also becomes mostly CP-odd and thus could have escaped detection at LEP. [Carena, Pilaftsis, Ellis, Wagner](#)

The CPX scenario

CPX Higgs

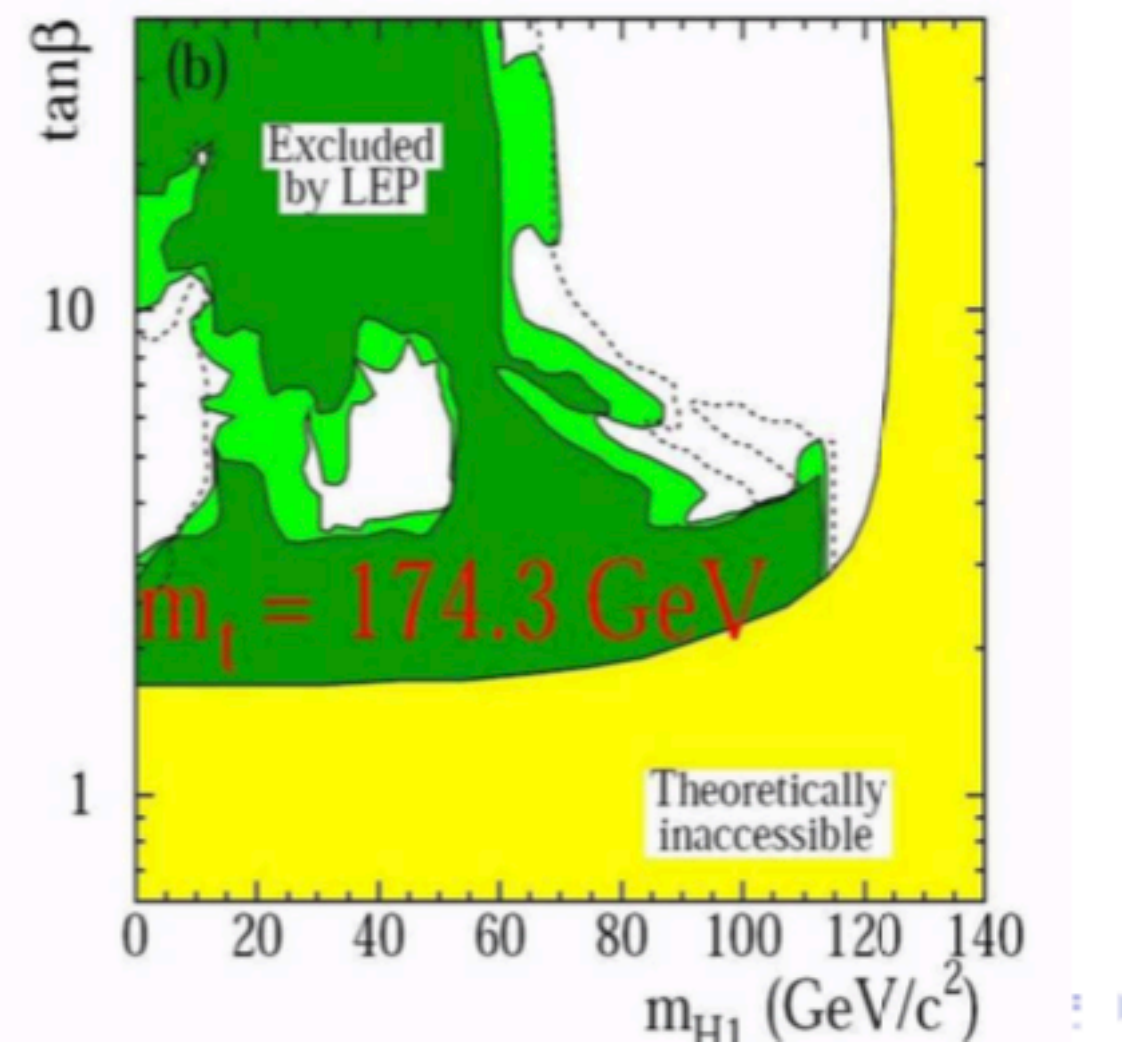
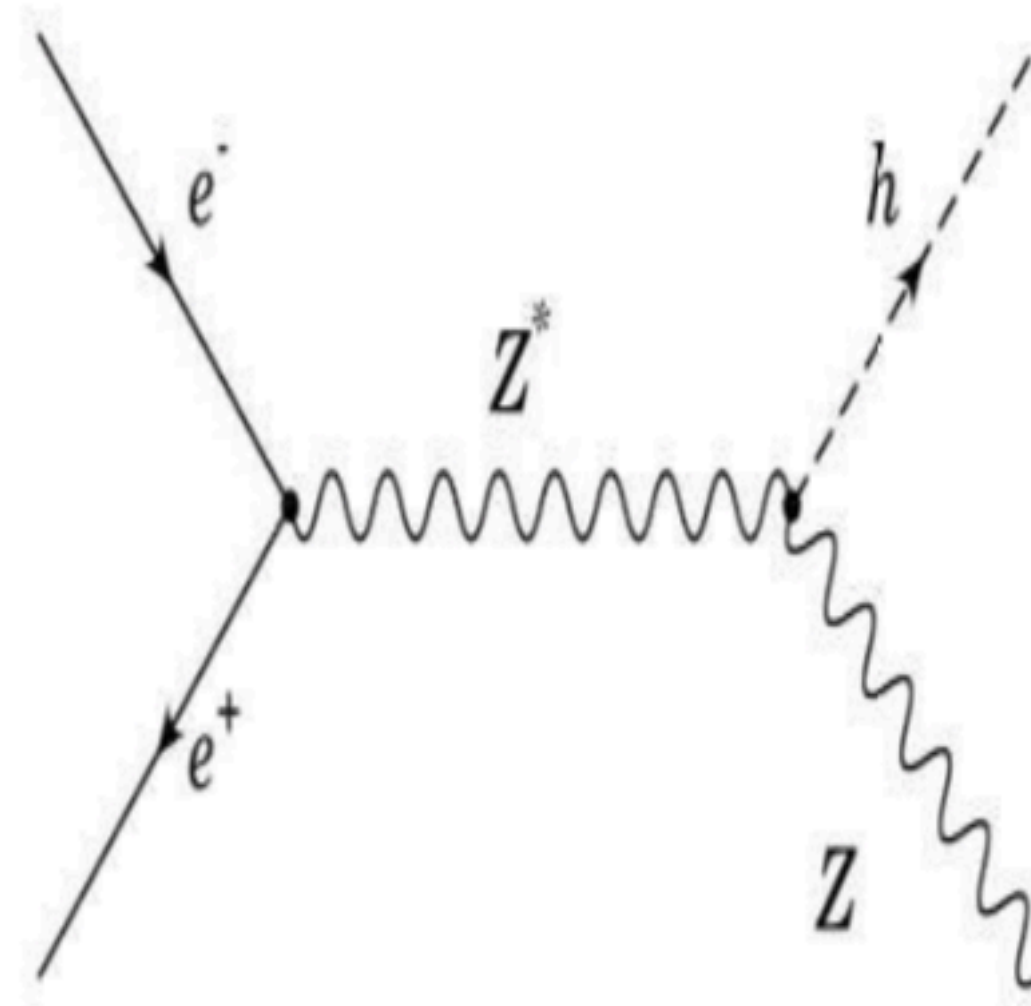
- A low lightest Higgs(h_1 , $m_{h_1} \leq 60$ GeV)
- $Z - Z - h_1$ coupling goes down.

Associate Higgs Production

$$\begin{aligned} pp &\rightarrow \tilde{t}_1 \tilde{t}_1^* h_1 \\ \tilde{t}_1 &\rightarrow b \chi_1^+ \rightarrow b W^+ \chi_1^0 \rightarrow b \ell^+ \nu_\ell \chi_1^0 \\ h_1 &\rightarrow b \bar{b} \end{aligned}$$

- **4-b partons + dilepton + missing p_T** can probe the 'hole' at LHC with $\mathcal{L}=30\text{fb}^{-1}$.
P.B, Aresh Krishna Datta, Amitava Datta, Biswarup Mukhopadhyaya; Phys. Rev. D 78, 015017 (2008)

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CP-violating Higgs production at the LHC

Higgs production in CPV-cascade decays

- Unlike MSSM, $H^\pm \rightarrow h_1 W^\pm$ mode is open here for light charged Higgs.
- Cascade decays of third generation squarks (\tilde{t}_1 and \tilde{b}_1) and gluino (\tilde{g}) could give signature at colliders.
- We produce at the LHC

$$pp \rightarrow \tilde{t}_1 \tilde{t}_1^*, \tilde{b}_1 \tilde{b}_1^*, \tilde{g} \tilde{g}$$

- Stop decays as:

$$\tilde{t}_1 \rightarrow t \chi_1^0 / b \chi_1^+ \rightarrow b H^+ / W^- \chi_1^0$$

- Similarly, sbottom decays as:

$$\tilde{b}_1 \rightarrow \tilde{t}_1 H^-$$

- A multi-channel analysis shows that, 5σ significance can be achieved with an integrated luminosity of $5-10 \text{ fb}^{-1}$.
P.Bandyopadhyay, JHEP08(2011)016

Two Higgses production

- Due to CP-mixing the heavier Higgses decay as: $h_{3,2} \rightarrow Z, h_1$.
- Unlike other processes Higgs pair production channels are clean
- We probe $pp \rightarrow h_1 h_{2,3}$ at the LHC.
 $\Rightarrow pp \rightarrow 4b + 2l$ signal topology.
P.B, Katri Huitu; arXiv:1106.5108 [hep-ph]
- This scenario is getting more constrained with at least one Higgs around 125 GeV, flavour and EMD constraints.
P.B, Katri Huitu; JHEP 1311 (2013) 058

Few comments

- One of Sakharov's condition: Departure from thermal equilibrium can be realized by the strong first order EWPT
- To have a first-order phase transition effective Higgs potential of the model should have several minima.
- Of special importance is the cubic term in effective potential

$$V_{eff}[\varphi, T] = D(T^2 - T_0^2)\varphi^2 - ET\varphi^3 + \frac{\lambda_T}{4}\varphi^4 + \dots,$$

- In the Standard Model the cubic term, ET^3 , is contributed only by the electroweak gauge bosons

$$v = 246 \text{ GeV} \implies E = \frac{2m_W^3 + m_Z^3}{4\pi v^3} \simeq 0.01$$

- Condition that the Higgs effective potential has two minima leads to the very small value for the Higgs self-coupling

$$\lambda \simeq 2E \simeq 0.02$$

- But $m_{h_{125}} = \sqrt{2\lambda}v = 125.5 \implies \lambda = 0.13$

- This means that within the Standard Model and MSSM the electroweak phase transition is a smooth crossover, i.e. of the second order

- For first order EWPT $m_{h_{SM}} \lesssim 73 \text{ GeV}$, $m_{h_{MSSM}} \lesssim 53 \text{ GeV}$

Conclusions

- So far we have observed one Higgs boson at 125 GeV
- More scalar is/are motivated from vacuum stability
- All standard and non-standard modes are yet to be explored
- Observation of Charged Higgs would be a direct proof of extended Higgs sector.
- Non-standard decay modes $h^\pm \rightarrow a_1 W^\pm$ and $h^\pm \rightarrow ZW^\pm$ are direct proofs of higher representations of Higgs sectors.
- Both SUSY and Non-SUSY Higgs extensions have their own merits
- Indirect searches can also give us some hints
- Electro-weak-ino sector has much more story to tell, so soft physics is important
- Finite temperature calculation along with CP-violation and Baryogenesis make such extensions very interesting

THANK

You!