

Supersymmetric Solutions of the Flavor Anomalies

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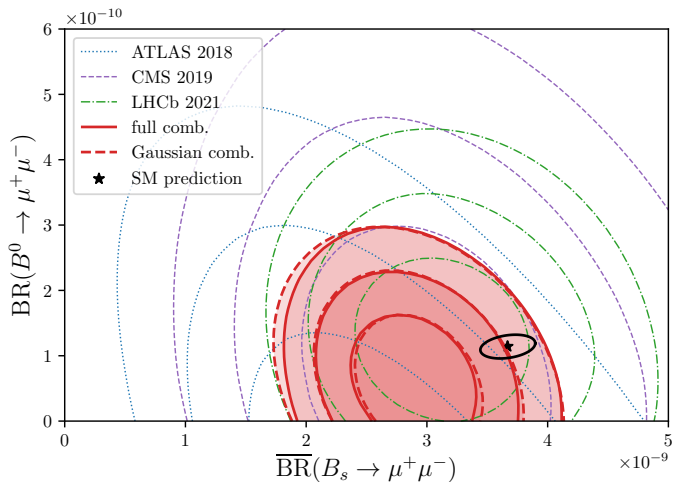


Anomalies 2021
November 10 - 12, 2021, Hyderabad

Overview of the Anomalies

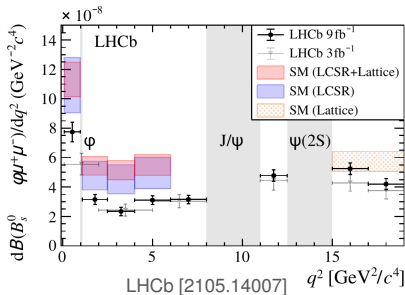
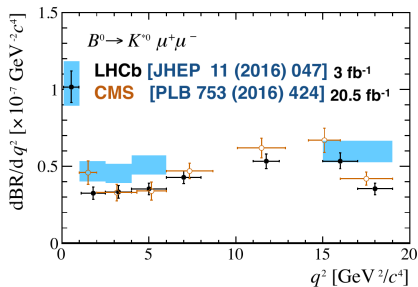
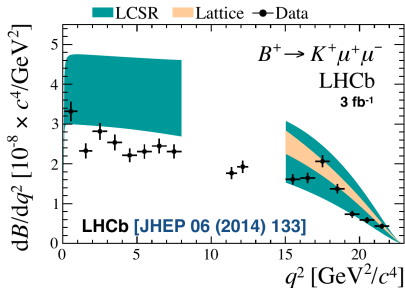
The $B_s \rightarrow \mu^+ \mu^-$ and $B_d \rightarrow \mu^+ \mu^-$ Decays

WA, Stangl 2103.13370; combination of LHCb 2108.09284, CMS 1910.12127, ATLAS 1812.03017



$\sim 2\sigma$ tension between SM and experiment

Semileptonic Branching Ratios



Experimental results for

$$\text{BR}(B \rightarrow K \mu \mu)$$

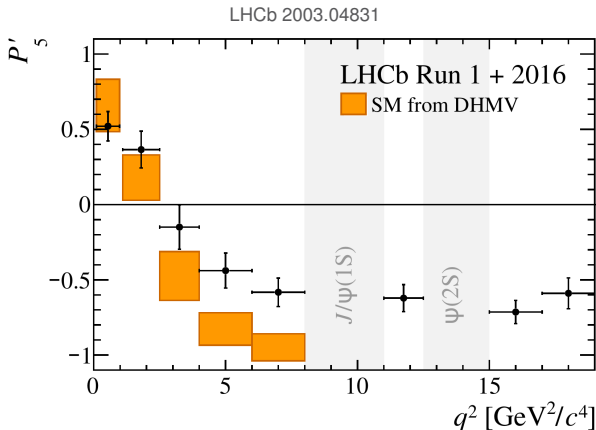
$$\text{BR}(B \rightarrow K^* \mu \mu)$$

$$\text{BR}(B_s \rightarrow \phi \mu \mu)$$

are consistently low
across many q^2 bins

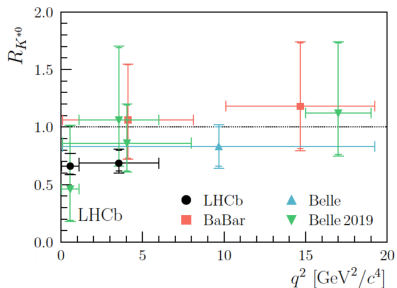
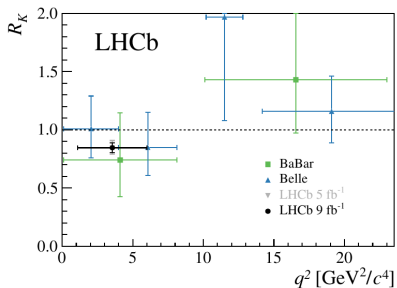
The P'_5 Anomaly

$P'_5 \sim$ a moment of the $B \rightarrow K^* \mu^+ \mu^-$ angular distribution



Anomaly persists in the latest update of $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ with 2016 data.
(Anomaly also seen in $B^\pm \rightarrow K^{*\pm} \mu^+ \mu^-$ LHCb 2012.13241)

Evidence for Lepton Flavor Universality Violation



$$R_{K^{(*)}} = \frac{BR(B \rightarrow K^{(*)} \mu \mu)}{BR(B \rightarrow K^{(*)} e e)}$$

$$R_{K^+}^{[1,6]} = 0.846^{+0.042+0.013}_{-0.039-0.012}$$

$$R_{K^{*0}}^{[0.045, 1.1]} = 0.66^{+0.11}_{-0.07} \pm 0.03$$

$$R_{K^{*0}}^{[1.1, 6]} = 0.69^{+0.11}_{-0.07} \pm 0.05$$

$$R_{K_S}^{[1.1, 6]} = 0.66^{+0.20+0.02}_{-0.14-0.04}$$

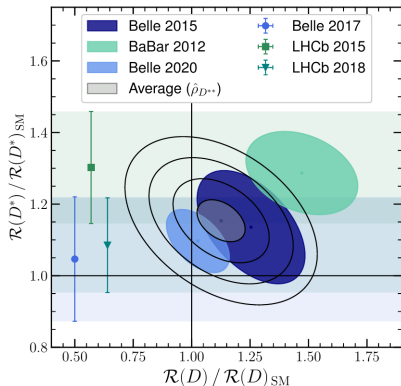
$$R_{K^{*+}}^{[0.045, 6]} = 0.70^{+0.18+0.03}_{-0.13-0.04}$$

$$R_{pK}^{[0.1, 6]} = 0.86^{+0.14}_{-0.11} \pm 0.05$$

LHCb 2103.11769

LFU in Charged Current Decays: R_D and R_{D^*}

Bernlochner, Franco Sevilla, Robinson, 2101.08326



$$R_D = \frac{BR(B \rightarrow D\tau\nu)}{BR(B \rightarrow D\ell\nu)}$$

$$R_{D^*} = \frac{BR(B \rightarrow D^*\tau\nu)}{BR(B \rightarrow D^*\ell\nu)}$$

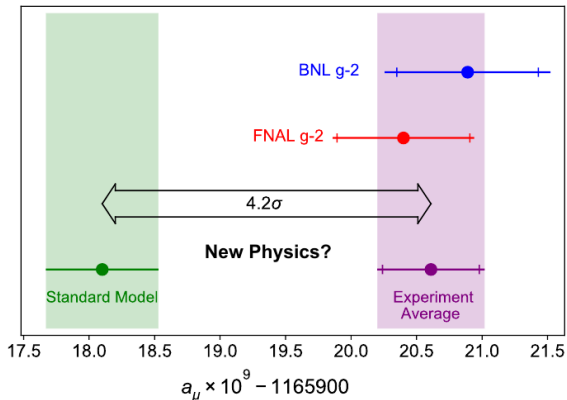
$$\begin{aligned} \ell = \mu, e & \quad (\text{BaBar/Belle}) \\ \ell = \mu & \quad (\text{LHCb}) \end{aligned}$$

$$R_D^{\text{exp}}/R_D^{\text{SM}} = 1.13 \pm 0.10, \quad R_{D^*}^{\text{exp}}/R_{D^*}^{\text{SM}} = 1.15 \pm 0.06$$

combined discrepancy with the SM: 3.6σ

(the heavy flavor averaging group quotes 3.1σ)

Anomalous Magnetic Moment of the Muon

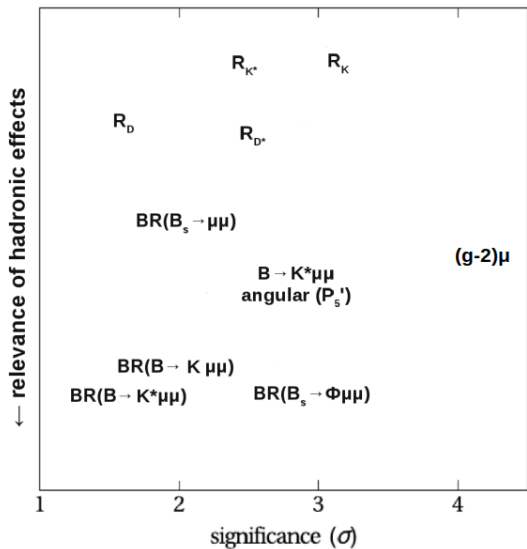


4.2 σ discrepancy between the experimental average (Fermilab g-2, 2104.03281) and the SM consensus (Aoyama et al. 2006.04822)

(see, however, the lattice results from BMW 2002.12347)

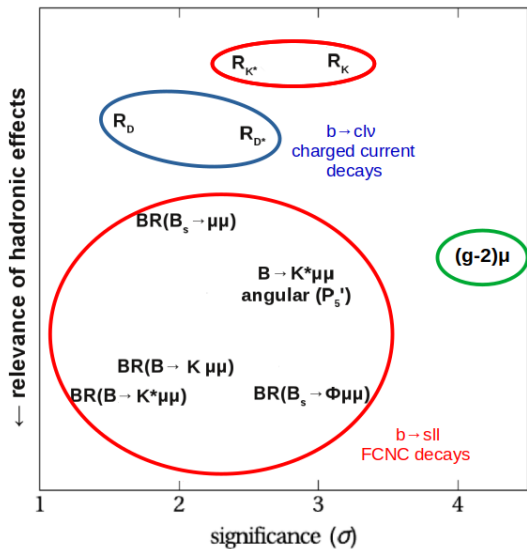
$$\Delta a_\mu = (251 \pm 59) \times 10^{-11}$$

(Selection of) Anomalies in 2021



(inspired by
Zoltan Ligeti)

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What Could It Be?

$B_s \rightarrow \mu\mu$
rate

semileptonic
rates

angular
observables

LFU
ratios

$(g - 2)_\mu$

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	$B_s \rightarrow \mu\mu$ rate	semileptonic rates	angular observables	LFU ratios	$(g-2)_\mu$
experimental issues?	?	?	?	?	?

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statistical fluctuations?	✓	✓	✓	✓	✗

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underestimated hadronic effects?	✗	✓	✓	✗	✓

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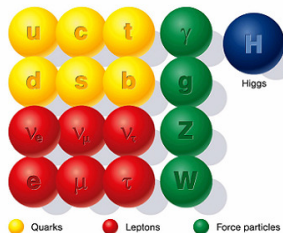
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experimental issues?	?	?	?	?	?
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parametric uncertainties?	✓	✓	✗	✗	✗
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New Physics?	✓	✓	✓	✓	✓

The Flavor Anomalies in the MSSM

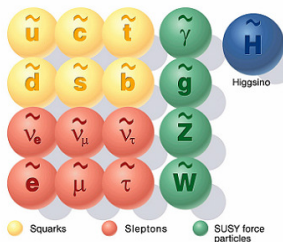
The Minimal Supersymmetric Standard Model

- Arguably **still the best motivated extension** of the Standard Model
- For a natural weak scale, need light Higgsinos, light stops, and relatively light gluinos
- First and **second generation of sfermions can be heavy** without spoiling the successes of the MSSM (naturalness, gauge coupling unification, dark matter, ...)

Standard particles

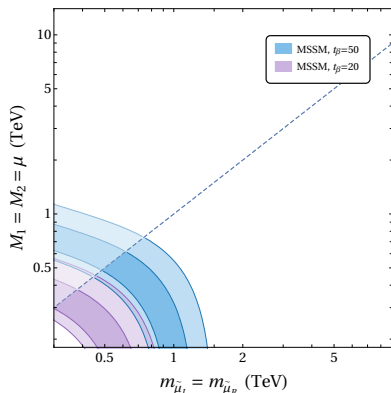


SUSY particles



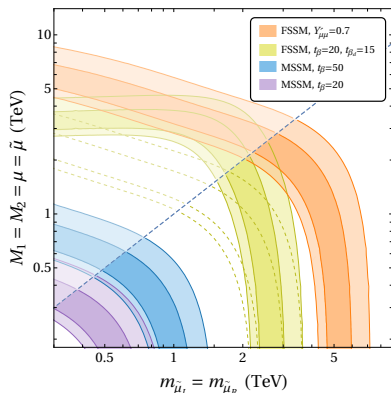
The Anomalous Magnetic Moment in the MSSM

- ▶ It is very well known that the **MSSM** can give sizeable contributions to $(g - 2)_\mu$ via **$\tan \beta$ enhanced** smuon chargino/neutralino loops
many many recent references
(apologies for the omission)
- ▶ Smuons, charginos, neutralinos need to be **pretty light**
- ▶ **Compressed spectra** to avoid existing LHC constraints
- ▶ Good discovery prospects at the high luminosity LHC and e^+e^- colliders (ILC, CLIC)



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- ▶ With **extended SUSY Higgs sectors**, smuons, charginos, neutralinos can be **significantly heavier**

WA, Gadam, Gori, Hamer 2104.08293

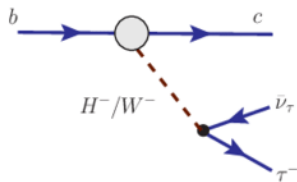
R_D and R_{D^*} in the MSSM

- ▶ There are tree level contributions to $B \rightarrow D^{(*)} \tau \nu$ from **charged Higgs exchange**

$$\frac{R_D}{R_D^{\text{SM}}} \sim 1 - 1.5 \frac{m_\tau m_b}{m_{H^\pm}^2} \tan^2 \beta$$

$$\frac{R_{D^*}}{R_{D^*}^{\text{SM}}} \sim 1 - 0.12 \frac{m_\tau m_b}{m_{H^\pm}^2} \tan^2 \beta$$

- ▶ Effect goes in the **wrong direction** and is much smaller for R_{D^*}



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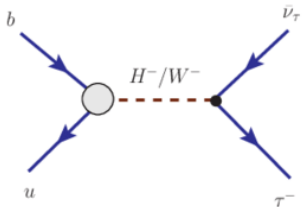
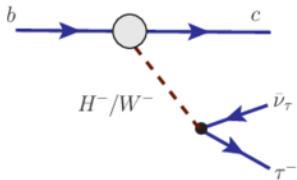
$$\frac{R_D}{R_D^{\text{SM}}} \sim 1 - 1.5 \frac{m_\tau m_b}{m_{H^\pm}^2} \tan^2 \beta$$

$$\frac{R_{D^*}}{R_{D^*}^{\text{SM}}} \sim 1 - 0.12 \frac{m_\tau m_b}{m_{H^\pm}^2} \tan^2 \beta$$

- ▶ Effect goes in the **wrong direction** and is much smaller for R_{D^*}
- ▶ Correlated with effect in $B \rightarrow \tau \nu$

$$\frac{\text{BR}(B \rightarrow \tau \nu)}{\text{BR}(B \rightarrow \tau \nu)_{\text{SM}}} \simeq \left(1 - \frac{m_B^2}{m_{H^\pm}^2} \tan^2 \beta \right)^2$$

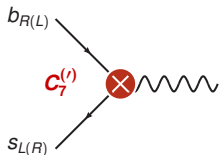
⇒ **Can't explain $R_{D^{(*)}}$ with charged Higgs exchange** in the MSSM



Interlude: R_K and R_{K^*} Model Independently

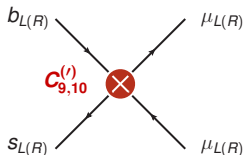
$$\mathcal{H}_{\text{eff}}^{b \rightarrow s} = -\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \frac{e^2}{16\pi^2} \sum_i (C_i \mathcal{O}_i + C'_i \mathcal{O}'_i)$$

magnetic dipole operators



$$C_7^{(i)} (\bar{s} \sigma_{\mu\nu} P_{R(L)} b) F^{\mu\nu}$$

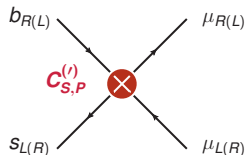
semileptonic operators



$$C_9^{(i)} (\bar{s} \gamma_\mu P_{L(R)} b) (\bar{\mu} \gamma^\mu \mu)$$

$$C_{10}^{(i)} (\bar{s} \gamma_\mu P_{L(R)} b) (\bar{\mu} \gamma^\mu \gamma_5 \mu)$$

scalar operators

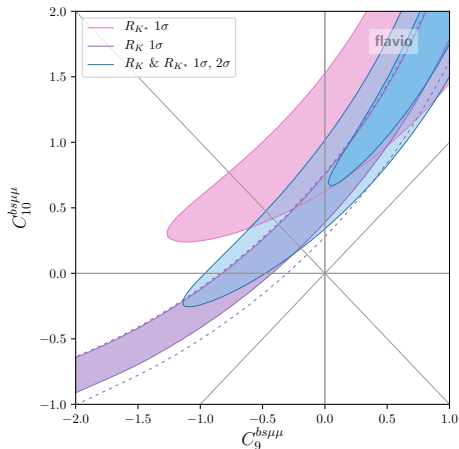


$$C_S^{(i)} (\bar{s} P_{R(L)} b) (\bar{\mu} P_{L(R)} \mu)$$

neglecting tensor operators and additional scalar operators

(they are dimension 8 in SMEFT: Alonso, Grinstein, Martin Camalich 1407.7044)

Interlude: Global Rare B Decay Fits



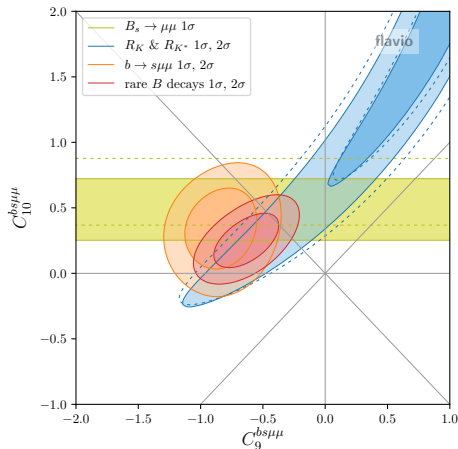
$$C_9^{bs\mu\mu}(\bar{s}\gamma_\alpha P_L b)(\bar{\mu}\gamma^\alpha \mu)$$

$$C_{10}^{bs\mu\mu}(\bar{s}\gamma_\alpha P_L b)(\bar{\mu}\gamma^\alpha \gamma_5 \mu)$$

- LFU ratios prefer non-standard C_{10} , but large degeneracy

WA, Stangl 2103.13370 (other recent fits: Geng et al. 2103.12738; Cornella et al. 2103.16558; Alguero et al. 2104.08921; Hurth et al. 2104.10058; Ciuchini et al. 2110.10126)

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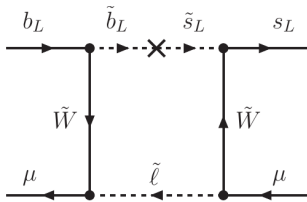
- ▶ LFU ratios prefer non-standard C_{10} , but large degeneracy
- ▶ $B_s \rightarrow \mu^+ \mu^-$ branching ratio shows slight preference for non-standard C_{10}
- ▶ $b \rightarrow s \mu \mu$ observables prefer non-standard C_9
- ▶ best fit point

$$C_9^{bs\mu\mu} \simeq -0.63$$

$$C_{10}^{bs\mu\mu} \simeq +0.25$$

WA, Stangl 2103.13370 (other recent fits: Geng et al. 2103.12738; Cornella et al. 2103.16558; Alguero et al. 2104.08921; Hurth et al. 2104.10058; Ciuchini et al. 2110.10126)

R_K and R_{K^*} in the MSSM



WA, Straub

1308.1501, 1411.3161

- only way to get lepton flavor non universal contribution to rare $b \rightarrow s l l$ decays is through box diagrams with light winos (or Binors) and large **non-universality in slepton masses**.

- requires an **extremely light spectrum** to get $C_9^{bs\mu\mu} \sim -0.5$:

winos and smuons around 100 GeV;

sbottoms around 500 GeV;

very challenging to hide this at the LHC...

The Flavor Anomalies and R-Parity Violation

The MSSM with R-Parity Violation

- give up on a dark matter candidate, but open up possibilities to address the flavor anomalies
- consider the lepton number violating *LQD* and *LLE* interactions (no baryon number violating *UDD* interactions to avoid constraints from proton decay)

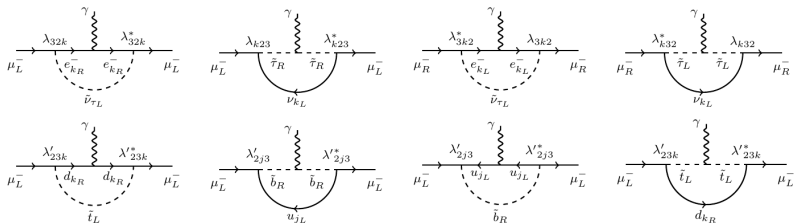
$$\mathcal{L}_{LQD} = \lambda'_{ijk} \left[\tilde{\nu}_{iL} \bar{d}_{kR} d_{jL} + \tilde{d}_{jL} \bar{d}_{kR} \nu_{iL} + \tilde{d}_{kR}^* \bar{\nu}_{iL}^c d_{jL} - \tilde{e}_{iL} \bar{d}_{kR} u_{jL} - \tilde{u}_{jL} \bar{d}_{kR} e_{iL} - \tilde{d}_{kR}^* \bar{e}_{iL}^c u_{jL} \right] + \text{H.c.}$$

$$\mathcal{L}_{LLE} = \frac{1}{2} \lambda_{ijk} \left[\tilde{\nu}_{iL} \bar{e}_{kR} e_{jL} + \tilde{e}_{jL} \bar{e}_{kR} \nu_{iL} + \tilde{e}_{kR}^* \bar{\nu}_{iL}^c e_{jL} - (i \leftrightarrow j) \right] + \text{H.c.}$$

- assume that **only the 3rd generation sfermions are light**
⇒ 7 λ couplings and 19 λ' couplings are relevant

→ **RPV3** (WA, Dev, Soni 1704.06659; WA, Dev, Soni, Sui 2002.12910; Dev, Soni, Xu 2106.15647)

The Anomalous Magnetic Moment with RPV3



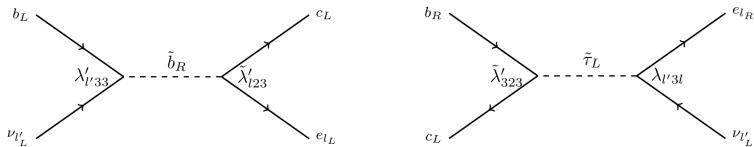
Kim, Kyae, Lee hep-ph/0103054

- 1-loop contributions from λ' and λ couplings (in addition to the standard MSSM contributions)

$$\Delta a_\mu = \frac{m_\mu^2}{96\pi^2} \sum_{k=1}^3 \left(\frac{2(|\lambda_{32k}|^2 + |\lambda_{3k2}|^2)}{m_{\tilde{\nu}_\tau}^2} - \frac{|\lambda_{3k2}|^2}{m_{\tilde{\tau}_L}^2} - \frac{|\lambda_{k23}|^2}{m_{\tilde{\tau}_R}^2} + \frac{3|\lambda'_{2k3}|^2}{m_{\tilde{b}_R}^2} \right)$$

- Need **light sbottoms and/or sneutrinos** with **large couplings** to get a relevant contribution in the right direction

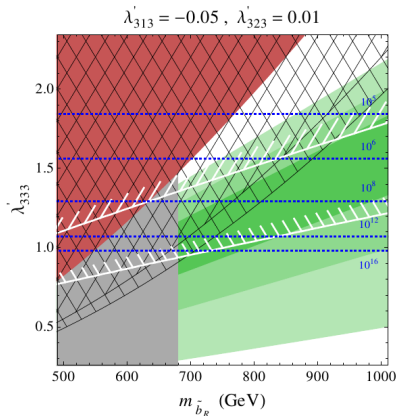
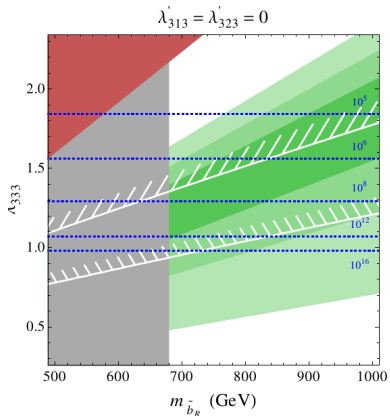
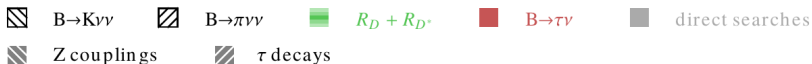
R_D and R_{D^*} with RPV3



Deshpande, He 1608.04817; WA, Dev, Soni 1704.06659; ...

- Tree level contributions from **sbottom or stau exchange**
- Stau behaves like a charged Higgs (but its couplings are less constrained). **Stau contribution disfavored** by $B_c \rightarrow \tau\nu$ branching ratio and kinematic distributions in $B \rightarrow D^{(*)}\tau\nu$.
- **Sbottom behaves like a leptoquark**. Chirality structure as preferred by model independent fits (Shi et al. 1905.08498; Murgui et al. 1904.09311; Asadi, Shih 1905.03311; Cheung et al. 2002.07272; ...)
- **Can address the $R_{D^{(*)}}$ anomalies** for sbottom masses $O(1 \text{ TeV})$ and couplings $\lambda' \sim O(1)$
- need to be careful to **keep $\mu - e$ universality** in $b \rightarrow c\ell\nu$

Viable Parameter Space

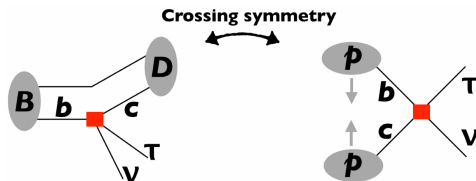


WA, Dev, Soni 1704.06659

Collider Signatures of $R_{D^{(*)}}$ Explanation

Expect non-standard
mono-tau production
at the LHC

(possibly in association
with b-jets)



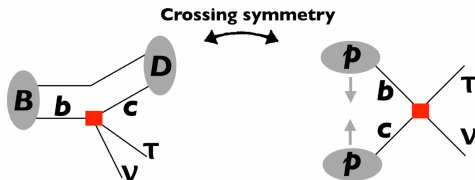
WA, Dev, Soni 1704.06659; Greljo et al. 1811.07920;

Marzocca et al. 2008.07541; ...

Collider Signatures of $R_{D^{(*)}}$ Explanation

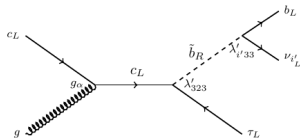
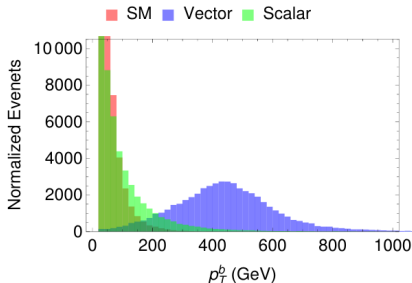
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WA, Dev, Soni 1704.06659; Greljo et al. 1811.07920;

Marzocca et al. 2008.07541; ...



► In RPV3, look for sbottom
production $gc \rightarrow \tilde{b}_T \rightarrow b\nu_T$

Implications for Neutrino Masses



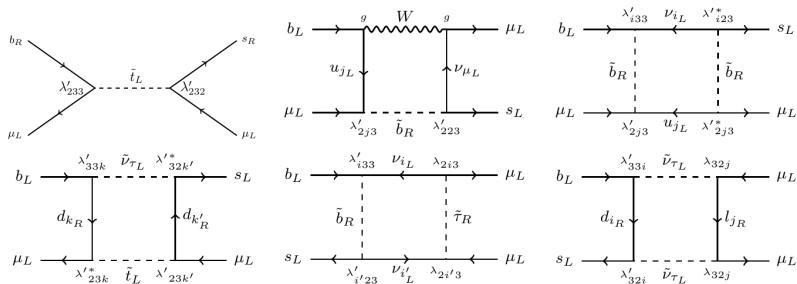
Barbier et al. hep-ph/0406039; WA, Dev, Soni 1704.06659

- The RPV couplings give also **1-loop contributions to Majorana neutrino masses**

$$(\hat{M}_\nu)_{ij} = (\hat{M}_\nu)_{ij}^{\text{tree}} + \frac{3}{8\pi^2} \frac{m_b^2 (A_b - \mu \tan \beta)}{m_{\tilde{b}}^2} \lambda'_{i33} \lambda'_{j33} + \frac{1}{8\pi^2} \frac{m_\tau^2 (A_\tau - \mu \tan \beta)}{m_{\tilde{\tau}}^2} \lambda_{i33} \lambda_{j33} + \dots$$

- Generic size of neutrino masses for sbottoms/staus masses of $O(1 \text{ TeV})$ and couplings of $O(1)$ is $\sim 0.1 \text{ MeV}$
- **Need cancellation** to obtain sub-eV neutrino masses

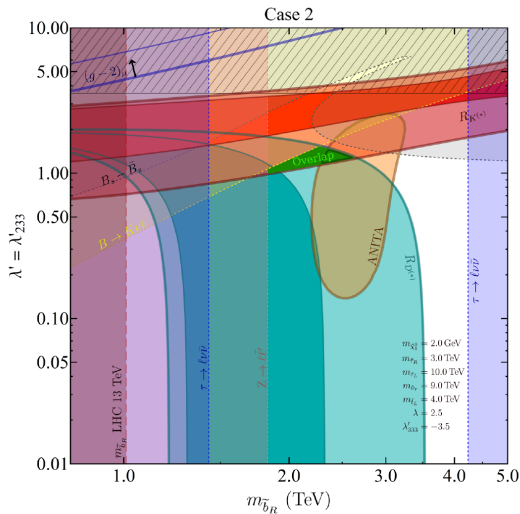
R_K and R_{K^*} with RPV



Das et al. 1705.09188; Earl Gregoire 1806.01343; Trifinopoulos 1807.01638; Hu, Huang 1912.03676;
 WA, Dev, Soni, Sui 2002.12910; Bardhan et al. 2107.10163

- Tree level contribution from stop exchange have the wrong chirality
- Several loop contributions with the right chirality and $C_9 = -C_{10}$
- Both λ and λ' couplings can be involved

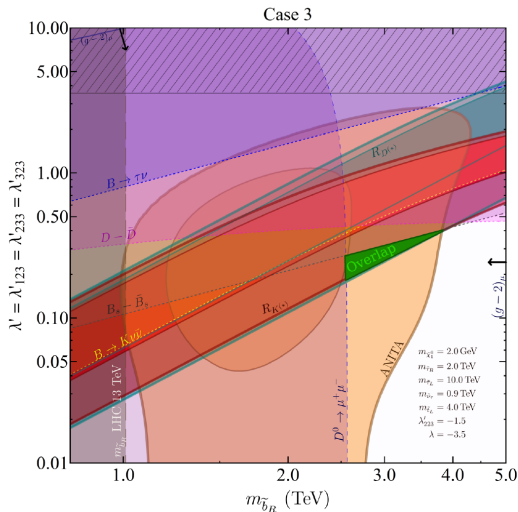
Combined Explanations of the Anomalies



WA, Dev, Soni, Sui 2002.12910

- We consider a few **benchmark scenarios**
- We include a very **long list of constraints**:
 - meson mixing;
 - rare decays;
 - Z decays;
 - lepton flavor violation;
 - direct LHC searches;
 - ...
- **Bonus**: can also explain ANITA events
Collins, Dev, Sui 1810.08479

Combined Explanations of the Anomalies

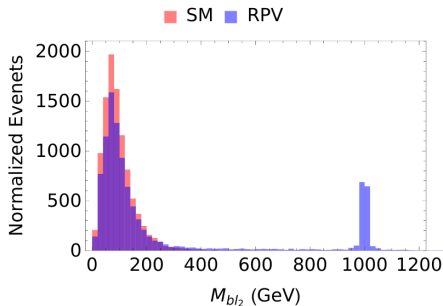
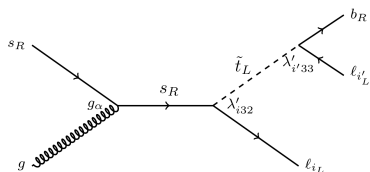


WA, Dev, Soni, Sui 2002.12910

- We consider a few benchmark scenarios
- We include a very long list of constraints:
 - meson mixing;
 - rare decays;
 - Z decays;
 - lepton flavor violation;
 - direct LHC searches;
 - ...
- **Bonus:** can also explain ANITA events
Collins, Dev, Sui 1810.08479

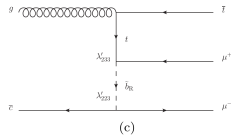
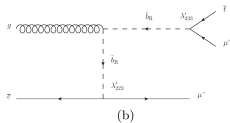
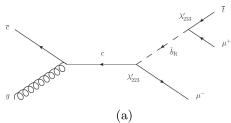
Collider Signatures of $R_{K^{(*)}}$ Explanation

- Based on **crossing symmetry** expect the processes $bs \rightarrow ll$, $gb \rightarrow sll$, and $gs \rightarrow bll$.
- In RPV3: for example **single stop production**, giving a $b\mu$ resonance.



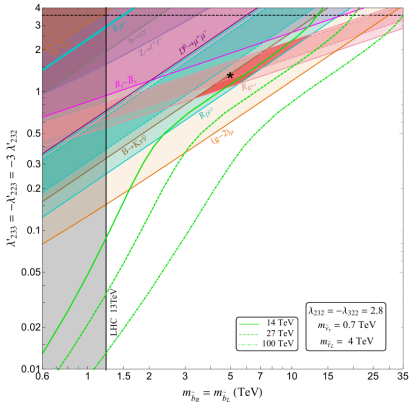
WA, Dev, Soni, Sui 2002.12910

More RPV3 Collider Signatures of the Anomalies

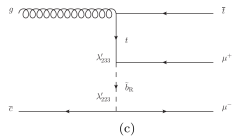
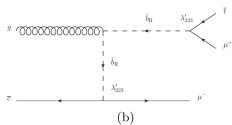
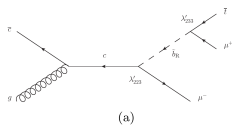


- $t\mu\mu$ production mediated by sbottoms

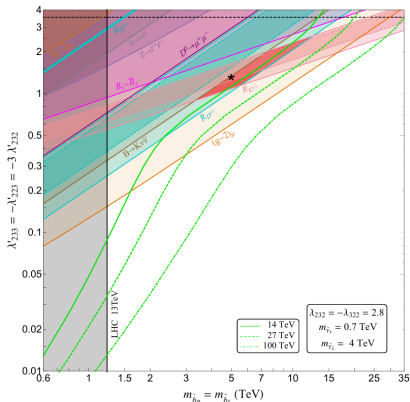
Dev, Soni, Xu 2106.15647



More RPV3 Collider Signatures of the Anomalies

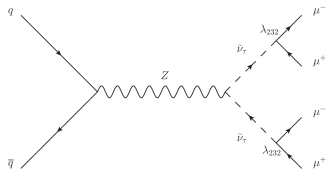


- $t\mu\mu$ production mediated by sbottoms



Dev, Soni, Xu 2106.15647

- pair of di-muon resonances



- ▶ Rare B decays and muon $g-2$ show persistent discrepancies with SM predictions.
- ▶ It's not possible to explain $R_{D^{(*)}}$ and $R_{K^{(*)}}$ in the MSSM.
- ▶ In the context of SUSY, need RPV interactions to explain hints for lepton flavor universality violation.
- ▶ In RPV3, combined explanations of the anomalies are strongly constrained but possible.
- ▶ RPV3 explanations lead to interesting collider signatures.