# Bounds on the Non-Oscillatory Explanations of MiniBooNE Excess

based on arXiv:2007.14411 in collaboration with O. Fischer and A. Yu. Smirnov

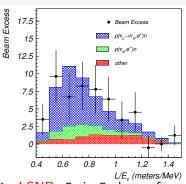
#### Vedran Brdar





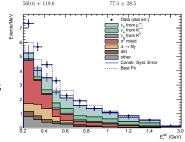
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#### LSND and MiniBooNE



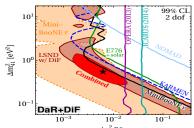
Process	Neutrino Mode	Antineutrino Mode
$\nu_{\mu} \& \bar{\nu}_{\mu} CCQE$	$107.6 \pm 28.2$	$12.9 \pm 4.3$
NC $\pi^0$	$732.3 \pm 95.5$	$112.3 \pm 11.5$
NC $\Delta \to N \gamma$	$251.9 \pm 35.2$	$34.7\pm5.4$
External Events	$109.8 \pm 15.9$	$15.3\pm2.8$
Other $\nu_{\mu}$ & $\bar{\nu}_{\mu}$	$130.8 \pm 33.4$	$22.3 \pm 3.5$
$\nu_e$ & $\bar{\nu}_e$ from $\mu^{\pm}$ Decay	$621.1 \pm 146.3$	$91.4 \pm 27.6$
$\nu_e$ & $\bar{\nu}_e$ from $K^{\pm}$ Decay	$280.7 \pm 61.2$	$51.2\pm11.0$
$\nu_e \& \bar{\nu}_e$ from $K_L^0$ Decay	$79.6 \pm 29.9$	$51.4\pm18.0$
Other $\nu_e$ & $\bar{\nu}_e$	$8.8 \pm 4.7$	$6.7\pm6.0$
Unconstrained Bkgd.	$2322.6 \pm 258.3$	$398.2 \pm 49.7$
Constrained Bkgd.	$2309.4 \pm 119.6$	$400.6\pm28.5$
Total Data	2870	478
Excess	$560.6 \pm 119.6$	$77.4 \pm 28.5$

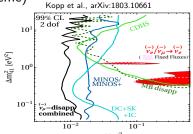
- ▶ LSND:  $\bar{\nu}_{\rm e}$  in  $\bar{\nu}_{\mu}$  beam from stopped pion source (> 3 $\sigma$ ) at  $L/E \sim 1 {\rm km \ GeV}^{-1}$  (arXiv:hep-ex/0104049)
  - MiniBooNE: reports electron-like event excess  $(4.8\sigma)$ ; in combination with LSND at  $6.1\sigma$  (arXiv:0812.2243, 1805.12028, 2006.16883)



#### eV-scale $\nu_s$ for LSND and MiniBooNE anomalies?

- Oscillation maxima for standard oscillations expected at
  - $\blacktriangleright$  L/E  $\sim$  500 km/GeV (from  $\Delta m_{31}^2 \sim 2.4 \times 10^{-3} \text{eV}^2$ )
  - $L/E \sim 15000 \text{ km/GeV (from } \Delta m_{21}^2 \sim 7.5 \times 10^{-5} \text{eV}^2)$
- ▶ the minimal solution for LSND and MiniBooNE requires an additional mass squared difference  $\Delta m_{41}^2 \sim 1 \, \text{eV}^2$ ; this calls for an introduction of eV-scale sterile neutrino (3+1 scheme)





while  $\nu_e$  appearance data supports eV-scale  $\nu_s$  explanation of LSND and MiniBooNE,  $\nu_\mu$  disappearance data puts such solution in strong tension and practically excludes this possibility  $\Rightarrow$  necessity for alternative models

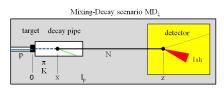
### From interactions of protons on target to 1 shower events

▶ The source of events are 8 GeV protons from Booster that hit the Beryllium target producing secondary particles. The 818 ton liquid scintillation detector observes the single shower events

$$p + A \text{ [target]} \rightarrow [X] \rightarrow 1 \text{sh events [detector]}$$

- The "black box", X, is assumed to be represented by a particle (or a system of particles) that are produced in the source  $(X_s)$  and evolve to detector where they interact or decay  $(X_d)$  producing 1 shower events
- X<sub>s</sub> can be produced
  - on target in pA collisions immediately
  - ▶ in decays (interactions) of known particles produced in the pA-collisions, such as  $\pi$ , K, heavy mesons. But those particles need to be charged!
  - From neutrinos  $u_{\mu}$  in detector or/and surrounding matter along the baseline
- $\triangleright X_d$ 
  - $ightharpoonup N o 
    u + \gamma, \ N o 
    u + e^+ + e^-$  (decay into particle(s)  $\xi$  that give shower)
  - $ightharpoonup N o 
    u + B, \quad B o e^+ + e^- \quad \text{or} \quad B o \gamma + \gamma$
  - $ightharpoonup N 
    ightharpoonup ... 
    u_e...$  followed by  $u_e$  scattering in the detector
  - $\triangleright$  N can also scatter  $\rightarrow$  additional smallness

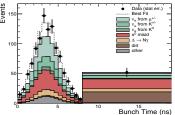
1)  $M_N D_{\xi}$ , Mixing - Decay scenario: the heavy neutrino N produced in the K and  $\pi$ -decay via mixing in  $\nu_{\mu}$  and decays as  $N \to N' + \xi$ 



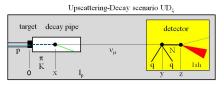
$$N_{\xi-s} = \epsilon A |U_{\mu 4}|^2 \int dE_N \frac{d\phi_N^0(E_N)}{dE_N} f_{\xi-s}(E_N) P_{dec}$$

 $P_{dec} pprox rac{d}{\lambda_N} e^{-I/\lambda_N}$ 

event excess peaks in the 8 ns window associated with beam bunch time, as expected from neutrino events in the detector  $\Longrightarrow m_N < 10 \, \text{MeV}$  for  $M_N D_\xi$  scenario



2)  $U_N D_\xi$ , Upscattering - decay scenario: N is produced in the  $\nu_\mu$  interactions with particles of medium between the source and the detector and in the detector. Then N decays in the detector, producing  $\xi$  state



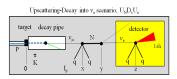
$$\frac{d\phi_N^{\sigma}(E_N)}{dE_N} \equiv \int dE_{\nu} \, \frac{d\phi_{\nu}(E_{\nu})}{dE_{\nu}} \, \frac{d\sigma(E_{\nu}, E_N)}{dE_N}$$

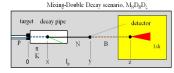
$$N_{\xi-s}^{in} = \epsilon V_d n_d \int dE_N f_{\xi-s}(E_N) \frac{d\phi_N^{\sigma}(E_{\nu})}{dE_N} \left[ 1 - \frac{\lambda_N}{d} (1 - e^{-d/\lambda_N}) \right]$$

- we also considered upscattering in the dirt as well as various detector subcomponents
- the models by Gninenko and Ballett et al. belong to this class of scenarios

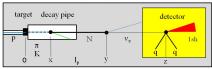
3)  $U_N D_B D_\xi$ , Upscattering - double decay scenario: N produced by  $\nu_\mu$  upscattering undergoes double decay:  $N \to B \to \xi$ . If B decays promptly, calculations match previous scenario

- 4)  $U_N D_{\nu} U_e$ , Upscattering-decay into  $\nu_e$  scenario: N produced by the  $\nu_{\mu}$  upscattering decays with emission of  $\nu_e$ , which then scatters in the detector via CCQE producing e shower
- 5)  $M_N D_B D_{\xi}$ , Mixing-double decay scenario: N produced via mixing decays invisibly into another new particle B, which in turn decays into (or with emission of)  $\xi$

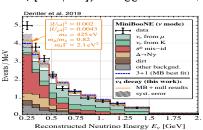




6)  $M_N D_{\nu} U_e$ , Mixing - Decay into  $\nu_e$  scenario: N is produced via mixing and decays with emission of  $\nu_e$ :  $N \to \nu_e + B$ . Then  $\nu_e$  upscatters in the detector, producing  $e^{\pm}$  Mixing-Decay  $\nu_e$  scenario MD, $U_e$ 



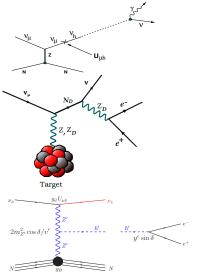
- ▶ for small N decay length  $c\tau^0 \to 0$  $N_{1e}^i \approx \sigma_{CC}^i V_i n_i B_N \phi_{\pi}^0 (1 - \text{Exp}[-I_T/\lambda_{\pi}]) \approx \sigma_{CC}^i V_i n_i B_N \phi_{\nu_{\mu}}$
- ▶ the spectrum for this scenario looks similar to the one in the 3+1 scenario
- $\triangleright$  viable N masses  $\mathcal{O}(\text{keV})$



# Non-oscillatory Explanations of MiniBooNE Anomaly

1 shower MiniBooNE events can be produced by e,  $\gamma$ , collimated  $e^+e^-$  pair and collimated  $\gamma\gamma$ 

- $M_N D_{\xi}$ : Fischer et al. (arXiv:1909.09561)
- U<sub>N</sub>D<sub>ξ</sub>:
   Gninenko (arXiv:0902.3802)
   Ballett et al. (arXiv:1808.02915)
- ► U<sub>N</sub>D<sub>B</sub>D<sub>ξ</sub>:
  Bertuzzo et al. (arXiv:1807.09877)
  Datta et al. (2005.08920)
  Dutta et al. (2006.01319)
  Abdallah et al. (2006.01948)
- $M_N D_{\nu} U_e$ :
  Dentler et al. (1911.01427)
  de Gouvea et al. (1911.01447)



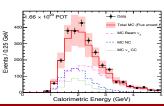
#### Strategy

- we employ several neutrino experiments to test aforementioned scenarios
- we normalize the numbers of events in a given detector, i, to the MiniBooNE excess

$$N_{\xi,exp}^{i} = N_{1sh,exp}^{MB} \frac{N_{\xi-s^{i}}^{i}}{N_{1sh}^{MB}} ,$$

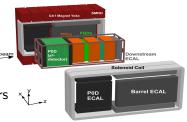
where  $N_{1sh,exp}^{MB} = 638$  and the remainder of the expression is the ratio of theoretical numbers of events in a given experiment and MiniBooNE

- in this way we ensure that a given scenario explains the MB excess; furthermore, various factors cancel in the ratio of predictions (mixing parameter, coupling constants...)
- for a given search, we use measured number of events as well as SM theory expectation and compute upper limit on the allowed number of new physics events at given confidence level



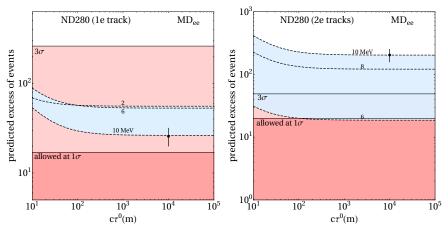
#### **Experiments**

- ► T2K ND280: sourced by 30 GeV protons that interact with a graphite target
- ▶ sub-detectors: the  $\pi^0$  detector P0D, the tracking detector containing the three Time Projection Chambers (TPC) filled by Ar gas, and two Fine Grained Detectors (FGD) filled by scintillatiors



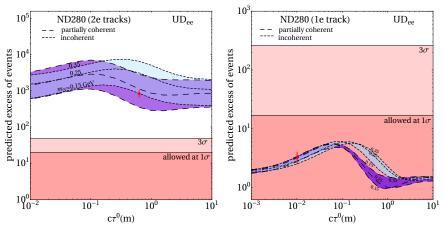
- ► MINER $\nu$ A: consists of scintillator strips; Good particle ID allows to identify 1e- from  $1\gamma$  and  $e^+e^-$  showers using the energy loss dE/dx
- ▶ PS191: was sourced by the PS proton beam with an energy of 19.2 GeV interacting with a beryllium target; measured excess of the e-like events in the calorimeter of  $23 \pm 8$  events
- ► NOνA: uses NuMI neutrino beam (120 GeV protons); composed of fine-grained cells of liquid scintillator

# $MD_{ee}$



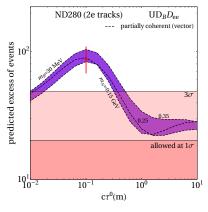
- $\triangleright$  timing limits impose consideration of  $\mathcal{O}(\text{MeV})$  masses
- ▶ invariant mass of  $e^+e^-$  pair used as a criterion for distinguishing between 1 and 2 showers



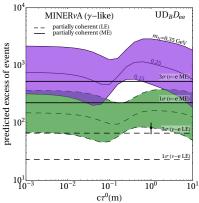


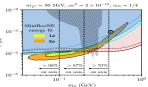
partially coherent cross section adopted from Bertuzzo et al. while incoherent one matches the benchmark point of Ballett et al.

# $UD_BD_{ee}$

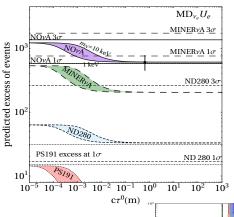


▶ Bertuzzo et al. model was tested by Arguelles et al. (arXiv:1812.08768) using MINER vA and CHARM-II data

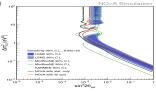




# $MD_{\nu_e}U_e$



NO $\nu$ A can test this scenario at  $2\sigma$  and disfavor MiniBooNE best fit point at  $\sim 3\sigma$ 



#### Summary

- a model independent study of the non-oscillatory explanations of the MiniBooNE excess was performed
- we carried out a systematic search of the simplest scenarios which can be classified by the number of new interaction points
- new physics scenarios allow to directly connect the observed MiniBooNE excess of events to expected excesses in other experiments (T2K ND280, MINERνA, PS-191, NOνA)
- each of the studied scenarios can be tested (portions of parameter space being already disfavored) using present neutrino data

# **BACKUP SLIDES**

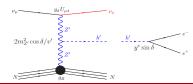
# Non-oscillatory Explanations of MiniBooNE Anomaly

1 shower MiniBooNE events can be produced by e,  $\gamma$ , collimated  $e^+e^-$  pair and collimated  $\gamma\gamma$ 

- ▶ Gninenko (arXiv:0902.3802): upscattering of  $\nu_{\mu}$  into  $\mathcal{O}(100)$  MeV right-handed neutrino which decays through magnetic moment operator into  $\gamma$   $\Longrightarrow |U_{\mu 4}|^2 \sim 10^{-3}$ ,  $\mu \sim 10^{-9} \mu_B$
- ▶ Bertuzzo et al. (arXiv:1807.09877): upscattering of  $\nu_{\mu}$  into  $\mathcal{O}(100)$  MeV right-handed neutrino which decays to on-shell Z' and  $\nu$ ; Z' then decays into collimated  $e^+e^-$  pair
- ▶ Ballett et al. (arXiv:1808.02915): identical particle content, however qualitatively different mechanism: right-handed neutrino N undergoes 3-body decay through the exchange of off-shell Z'; final state e<sup>+</sup>e<sup>-</sup> pair can mimic MiniBooNE 1 shower signature (this realization typically yields longer N lifetimes in comparison to previous scenario with on-shell Z')

# Non-oscillatory Explanations of MiniBooNE Anomaly II

- Fischer et al. (arXiv:1909.09561): production of N via mixing in  $\nu_{\mu}$  in the decay pipe and further radiative decay along the beamline and mainly in the detector
- Dentler et al. (1911.01427) and de Gouvea et al. (1911.01447): N production via mixing in the decay pipe and then decay  $N \to \nu_e \phi$  along the baseline with emission of  $\nu_e$  which then produces electron via CCQE scattering in the detector ( $\phi$  decay into  $\nu_e \bar{\nu}_e$  also contributes to the signal)
- ▶ Datta et al. (2005.08920) and Dutta et al. (2006.01319): analogous realization to Bertuzzo et al.; the crucial difference is usage of scalar  $\phi$  instead of Z' for the purpose of relaxing MINER $\nu$ A limit
- Abdallah et al. (2006.01948): production of the light scalar B in upscattering of  $\nu_{\mu}$ , which then decays as  $B \rightarrow e^+e^-$

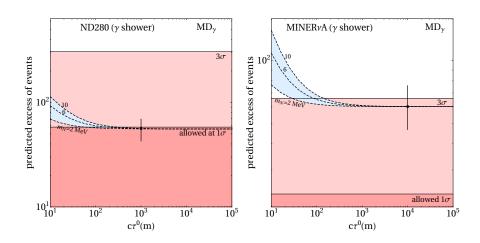


# Experiments and Searches

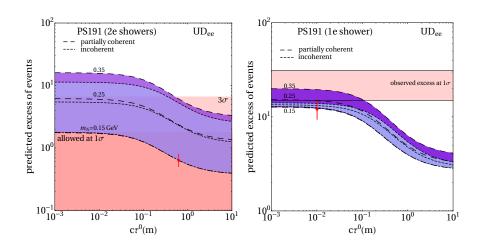
experiment	MiniBooNE	T2K	NOMAD	PS191	$MINER \nu A$	$NO\nu A$
area (m <sup>2</sup> )	$36\pi$	3.47	6.76	18	1.71	12.39
$\epsilon$	0.1	0.3	0.08	0.7	0.73	0.65
d (m)	$2/3 \cdot 12$	$d_1 = 1$ , $d_2 = 0.9$	3.7	3.55	3	8
$l_p$ (m)	50	94	290	49.1	675	675
POT $(\nu + \bar{\nu} \text{ mode})$	$3 \times 10^{21}$	$1.821 \times 10^{21}$	$2.2 \times 10^{19}$	$0.86 \times 10^{19}$	$3.43 \times 10^{20}$	$1.66 \times \cdot 10^{20}$
M (tonnes)	818	$m_{P0D} = 15.8, m = 1.1$	112	20	6.1	300
$\nu$ energy range (GeV)	[0.1 - 5]	[0.1 - 10]	[5 - 200]	[0.1, 5]	[0.1 - 20]	[0.1 - 20]

Experiment	Analysis	Signature	Upper limit $1\sigma/3\sigma$	
T2K ND280	Heavy neutrino decays	$e^+e^-$	20/49	
	CCQE electrons	$e^{-} (e^{+})$	17/261	
	CCQE electrons	single $\gamma$	58/305	
$NO\nu A$	CCQE electrons	$e^-$	577/1355	
$MINER\nu A$	diffractive $\pi^0$ production	$\gamma$	211/632	
	CCQE electrons	$e^{-} (e^{+})$	757/1725	
	Neutrino electron scattering	EM shower, or $\gamma, ee$	23/66	
	Neutrino electron scattering	EM shower, or $\gamma, ee$	223/526	
NOMAD	Single photon search	single $\gamma$	18/50	
PS191	Heavy neutrino decays	displaced vertex	1.84/6.61	
	Neutrino oscillation	electron-like events	$23 \pm 8$	

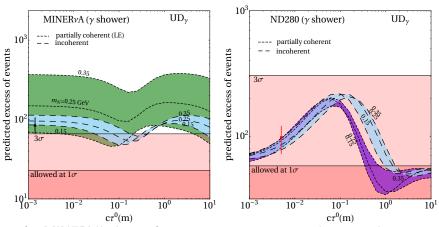
# $MD_{\gamma}$











- lacktriangle for MINERVA, limits from u-e scattering search are used
- ightharpoonup important cut on  $E\theta^2$